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A systematic optical study of phonon properties in optimally doped $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ single crystals

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Abstract

A systematic optical study of infrared-active phonons is carried out for optimally doped $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ (Bi2212) single crystals with E -vector along c -axis as well as in the ab -plane at 295 K. Strong phonon anisotropy has been observed in our ab -plane conductivity measurements. Ten infrared-active phonons have been observed along c -axis. Many of these c -axis infrared-active phonons have energies that are very close to the Raman-active A_{1g} modes. The group theory analysis based on the tetragonal group (D_{4h}) is inadequate to explain our experimental observations. © 2002 Elsevier Science B.V. All rights reserved.

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1. Introduction

In view of the renewed suggestion that strong electron-phonon interaction [1] is the mechanism for superconductivity in high T_c cuprates, phonon properties need to be examined more carefully in these systems. Phonons in high T_c cuprates has been the focus of many Raman works [2]. The vibrational spectra of high T_c cuprates are difficult to interpret and this is particularly true in the case of $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta}$ (Bi2212) with the presence of a structural modulation [3]. Infrared spectroscopy might shed some light on these questions. Using the vibrational spectra derived from both the infrared and Raman spectroscopy, a more com-

plete understanding of phonon properties of Bi2212 can be reached. Group theory analysis based on a tetragonal cell (space group D_{4h}) in Bi2212 [4], gives 30 modes at the center of the Brillouin zone for Bi2212 [5]: $\Gamma \equiv 6A_{1g} + B_{1g} + 7E_g + 7A_{2u} + 8E_u + B_{2u}$. Of the total of 30 modes: two are acoustic ($A_{2u} + E_u$) and therefore not infrared or Raman-active; thirteen ($6A_{2u} + 7E_u$) are infrared-active; fourteen ($6A_{1g} + B_{1g} + 7E_g$) are Raman-active and one is silent (B_{2u}).

The ab -plane optical conductivity of optimally doped Bi2212 single crystals has been measured extensively [6,7], but no lattice modes have been observed directly. However, only one brief study on the c -axis conductivity in Bi2212 was reported [8]. In this paper, a systematic optical study of infrared-active phonons is carried out for optimally doped Bi2212 single crystals with E -vector

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along c -axis as well as in the ab -plane at 295 K. Strong phonon anisotropy has been observed in our ab -plane conductivity measurements. Ten infrared-active phonons have been identified along c -axis. Many of these c -axis infrared-active phonons have energies that are very close to the Raman-active A_{1g} modes. The group theory analysis based on the tetragonal group (D_{4h}) is inadequate to explain our experimental observations.

2. Experimental details

For this study, large optimally doped Bi2212 single crystals are grown using traveling-surface-floating-zone (TSFZ) method [9]. The Bi2212 single crystals are mounted on an optically black cone, and the reflectance is measured in a near-normal-incidence arrangement from $\sim 50 \text{ cm}^{-1}$ to over $20,000 \text{ cm}^{-1}$ on a Bruker IFS 66v/S. The resolution of this infrared work is 2 cm^{-1} in the phonon region. The absolute reflectance is determined by evaporating a gold film ($\sim 100 \text{ nm}$ in thickness) in situ in an ultra-high vacuum ($\sim 1 \times 10^{-8} \text{ Torr}$). This comparison to the gold reflectivity provides an absolute reflectivity scale. The details of this technique are described elsewhere [10].

3. Experimental results and discussion

In Fig. 1, reflectance spectra are shown for optimally doped Bi2212 single crystals with E -vector parallel to a , b and c -axes at 295 K. In Fig. 1(a) and (b), the low frequency ab -plane reflectance spectra are given. Because of the unprecedented high signal-to-noise ratio (~ 1000) achieved using our in situ evaporation technique, structures due to infrared-active ab -plane phonons have been observed for the first time in optical measurements of optimally doped Bi2212. The reflectance for $E||b$ -axis is about 0.3% lower than that of for $E||a$ -axis in this frequency region consistent with the previous observations [7]. The reflectance spectrum for $E||c$ -axis is shown in

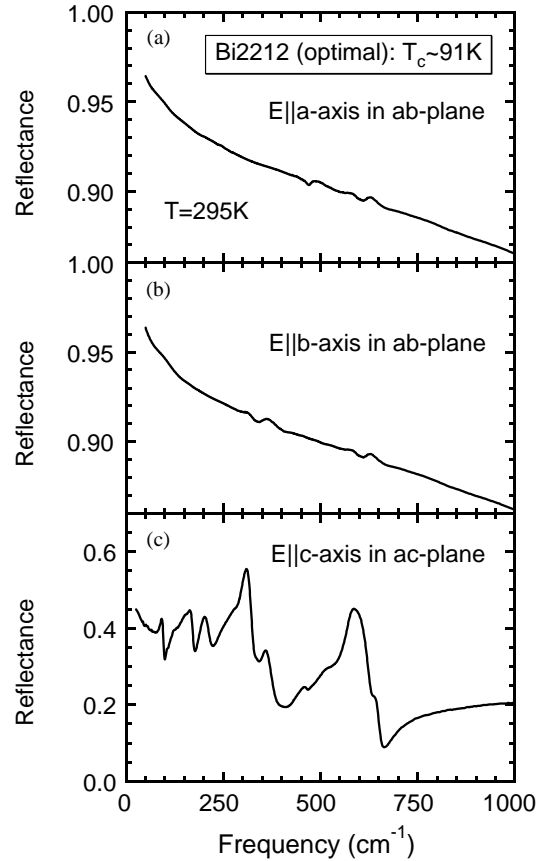


Fig. 1. The measured low frequency reflectance of optimally doped Bi2212 single crystals with the E -vector parallel to a , b and c -axes at 295 K. (a) The reflectance with $E||a$ -axis in ab -plane; (b) the reflectance with $E||b$ -axis in ab -plane; and (c) the reflectance with $E||c$ -axis in ac -plane.

Fig. 1(c). Many phonon features can be observed in the c -axis reflectance spectrum.

We have measured the reflectance over a wide region $\sim 50 \text{ cm}^{-1}$ to over $20,000 \text{ cm}^{-1}$ at 295 K to carry out a reliable Kramers-Kronig analysis. The results of such an analysis are shown in Fig. 2. Two ab -plane infrared-active TO phonons can be identified at 473.2 and 612.8 cm^{-1} with $E||a$ -axis marked as A and B, respectively, in σ_1 shown in Fig. 2(a). While the 612.8 cm^{-1} mode is also observed with $E||b$ -axis, the 473.2 cm^{-1} mode is completely absent and a new mode appears at 325.2 cm^{-1} marked as C in Fig. 2(b). It is perhaps not surprising that only three ab -plane

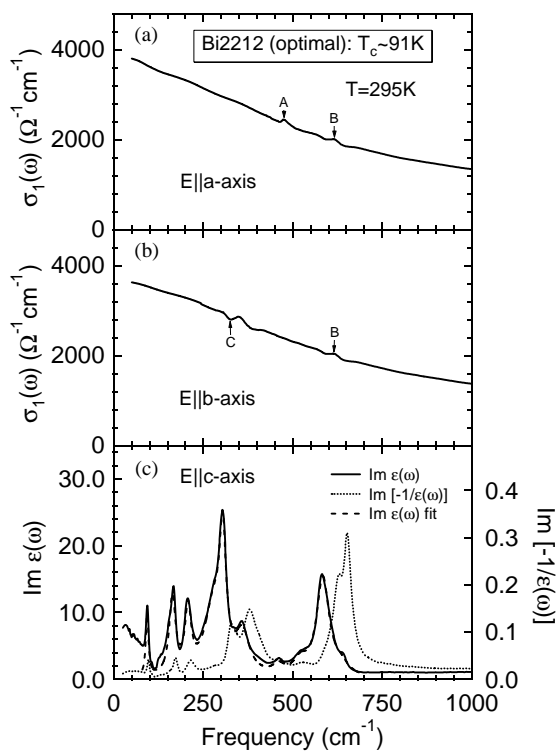


Fig. 2. Optical phonons in optimally doped Bi2212 single crystals at 295 K. (a) Two optical phonons at 473.2 and 612.8 cm^{-1} are observed in σ_1 with $E||a$ -axis in ab -plane. (b) Two optical phonons at 325.2 and 612.8 cm^{-1} are observed in σ_1 with $E||b$ -axis in ab -plane. (c) Ten optical TO phonons observed in $\text{Im } \epsilon(\omega)$ with $E||c$ -axis in ac -plane with a fit with ten Lorentzians; ten corresponding optical LO phonons are observed in the loss function $\text{Im}[-1/\epsilon(\omega)]$.

infrared-active modes are observed in σ_1 instead of $7E_u$ modes predicted by the D_{4h} space group. Many of these in-plane modes would be screened by the charge carriers in the ab -plane particularly those involving the vibrations within the CuO_2 planes. Therefore these three infrared-active in-plane modes at 325.2, 473.2 and 612.8 cm^{-1} are quite likely originated from in-plane oxygen vibrations in the BiO and SrO layers. The c -axis infrared-active phonons are shown in Fig. 2(c) at 295 K. It is better to examine the imaginary part of the dielectric function directly instead of plotting σ_1 for $E||c$ -axis. Ten c -axis infrared-active TO phonons can be identified in $\text{Im } \epsilon(\omega)$ with the

corresponding LO phonons derived from the loss function $\text{Im}[-1/\epsilon(\omega)]$ given in bracket: 93.3 (96.4), 164.1 (171.6), 207.7 (213.4), 280.5 (283.5), 304.4 (327.8), 360.3 (378.6), 460.2 (463.5), 521.4 (527.1), 584.4 (628.7) and 625.1 (651.8) cm^{-1} . More c -axis infrared-active phonons are observed in the present work than the $6A_{2u}$ phonons predicted by the D_{4h} space group.

The observed strong ab -plane phonon anisotropy is quite surprising giving the small orthorhombic distortions in optimally doped Bi2212. The strong polarization dependence of the 325.2 and 473.2 cm^{-1} modes indicates that they might be quasi-one-dimensional oxygen vibrations along b - and a -axis, respectively, ab -plane phonon anisotropy in optimally doped Bi2212 is inconsistent with the simplified D_{4h} space group assignment. The other indication of the inadequacy of the D_{4h} space group assignment is the fact that ten infrared-active c -axis phonons have been observed in this work while the D_{4h} space group predicts only six. In fact, our conclusion is similar to the Raman study on Bi2212 by Liu et al. [3] where they have demonstrated that the large number of peaks of A_{1g} symmetry observed experimentally is inconsistent with the D_{4h} space group and have used an alternative analysis based on an orthorhombic cell to explain their data. However, Liu et al. [3] have chosen a centrosymmetric space group $Bbmb$ [11] to describe the orthorhombic subcell. However, many of the infrared-active phonons have energies that are very close to that of the Raman-active modes measured by Liu et al. [3]. Based on this observation, we conclude that it is more reasonable to use the space group without centrosymmetry, i.e., $A2aa$ [12] for Bi2212. A detailed group theory analysis is needed to account for all the observed infrared-active and Raman-active phonons in Bi2212.

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